

Dr. Siegf. Guggenheimer

ELECTROMAGNETIC THEORY.

BY

OLIVER HEAVISIDE.



VOLUME I.

LONDON :
"THE ELECTRICIAN" PRINTING AND PUBLISHING COMPANY,
LIMITED.

SALISBURY COURT, FLEET STREET, E.C.

1893.

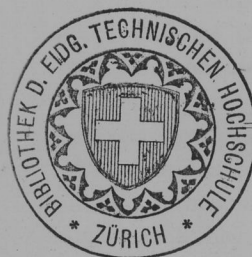
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to suit the case of inductive leakage, whilst R and L remain as before, the effective resistance and inductance per unit length of line. The solutions should be examined through a wide range of frequency in order to arrive at a good knowledge of how the circuit would behave to telephonic currents, with different values of λ and ρ . It is, of course, necessary that ρ should not be so small as to produce too great a loss in transit.

Notice that λ and ρ are immensely big when the unit length is a centimetre. Remember that m equal coils in parallel behave like one coil of one m^{th} the resistance and inductance of the individual coils. Thus, when ten shunts at the rate of one per kilometre take the place of one shunt per ten kilometres, each of the ten should have ten times the resistance and inductance of the one. This leads to physically monstrous results when we do as was done above, and make the action uniform, as it implies infinitely large inductance of an infinitely small section of the leakance. But that need not cause any alarm when we are employing the ideal case for calculating purposes. What is more important is that a stop would soon be set to the multiplication of the shunts.

At the particular frequency making the effective permittance vanish, the solution may be obtained simply from the fundamental formulæ, since Y reduces to a constant, viz., the value of K in (25), S being zero; and only the rationalisation of Z has to be attended to. If, further, we assume that $L=0$, then Z is also constant, viz., R , and the solutions are simply

$$V = e^{-\alpha(RK)^{\frac{1}{2}}V_0} \dots \dots \dots (26)$$

$$C = \left(\frac{K}{R}\right)^{\frac{1}{2}} V \dots \dots \dots (27)$$

These represent a very remarkable state of affairs throughout the cable. But it is thoroughly deceptive, because even if the inductance of the line were quite unimportant under normal conditions, the effective cancellation of the permittance would make it important under the present circumstances; that is, L cannot be neglected. Including it, when $S=0$, the value of P reduces to

$$P = \left(\frac{1}{2}\right)^{\frac{1}{2}} \left(\frac{\rho}{\rho^2 + \lambda^2 n^2}\right)^{\frac{1}{2}} \{R + (R^2 + L^2 n^2)^{\frac{1}{2}}\}, \quad (28)$$

and the current amplitude, according to (21), becomes

$$(C) = \frac{\rho^{\frac{1}{2}} e^{-P\sigma}}{(\rho^2 + \lambda^2 n^2)^{\frac{1}{2}} (R^2 + L^2 n^2)^{\frac{1}{2}}} \dots \dots \dots (29)$$

in which, of course, ρ and λn are not independent, having a relation fixed between them by the vanishing of S in (25). By giving a suitable value to σ , for instance, $\frac{1}{4}$ microfarad per kilometre (which is a suitable unit of length to employ), and also fixing the frequency, we may readily apply the formulæ to estimate the attenuation with different values of the insulation resistance or its equivalent.

APPENDIX B.

A GRAVITATIONAL AND ELECTROMAGNETIC ANALOGY.

Part I.

To form any notion at all of the flux of gravitational energy, we must first localise the energy. In this respect it resembles the legendary hare in the cookery book. Whether the notion will turn out to be a useful one is a matter for subsequent discovery. For this, also, there is a well-known gastronomical analogy.

Now, bearing in mind the successful manner in which Maxwell's localisation of electric and magnetic energy in his ether lends itself to theoretical reasoning, the suggestion is very natural that we should attempt to localise gravitational energy in a similar manner, its density to depend upon the square of the intensity of the force, especially because the law of the inverse squares is involved throughout.

Certain portions of space are supposed to be occupied by matter, and its amount is supposed to be invariable. Furthermore, it is assumed to have personal identity, so that the position and motion of a definite particle of matter are definite, at any rate relative to an assumed fixed space. Matter is recog-

nised by the property of inertia, whereby it tends to persist in the state of motion it possesses; and any change in the motion is ascribed to the action of force, of which the proper measure is, therefore, the rate of change of quantity of motion, or momentum.

Let ρ be the density of matter, and e the intensity of force, or the force per unit matter, then

$$\mathbf{F} = e\rho \quad \dots \quad (1)$$

expresses the moving force on ρ , which has its equivalent in increase of the momentum. There are so many forces nowadays of a generalised nature, that perhaps the expression "moving force" may be permitted for distinctness, although it may have been formerly abused and afterwards tabooed.

Now the force \mathbf{F} , or the intensity e , may have many origins, but the only one we are concerned with here is the gravitational force. This appears to depend solely upon the distribution of the matter, independently of other circumstances, and its operation is concisely expressed by Newton's law, that there is a mutual attraction between any two particles of matter, which varies as the product of their masses and inversely as the square of their distance. Let e now be the intensity of gravitational force, and \mathbf{F} the resultant moving force, due to all the matter. Then e is the space-variation of a potential, say,

$$e = \nabla P, \quad \dots \quad (2)$$

and the potential is found from the distribution of matter by*

$$P = \text{pot} \frac{\rho}{c} = \sum \frac{\rho}{4\pi cr},$$

where c is a constant. This implies that the speed of propagation of the gravitative influence is infinitely great.†

Now when matter is allowed to fall together from any configuration to a closer one, the work done by the gravitational force is expressed by the increase made in the quantity $\sum \frac{1}{2} P\rho$. This is identically the same as the quantity $\sum \frac{1}{2} ce^2$ summed through all space. If, for example, the matter be given

* See § 133.

† The density is expressed in terms of the intensity of force by $\rho = \text{conv} e$. This is also the case in the extended analogy later, when the lines of e are slightly shifted.

initially in a state of infinitely fine division, infinitely widely separated, then the work done by the gravitational force in passing to any other configuration is $\sum \frac{1}{2} P\rho$ or $\sum \frac{1}{2} ce^2$, which therefore expresses the "exhaustion of potential energy."* We may therefore assume that $\frac{1}{2} ce^2$ expresses the exhaustion of potential energy per unit volume of the medium. The equivalent of the exhaustion of potential energy is, of course, the gain of kinetic energy, if no other forces have been in action.

We can now express the flux of energy. We may compare the present problem with that of the motion of electrification. If moved about slowly in a dielectric, the electric force is appreciably the static distribution. Nevertheless, the flux of energy depends upon the magnetic force as well. It may, indeed, be represented in another way, without introducing the magnetic force, but then the formula would not be sufficiently comprehensive to suit other cases. Now what is there analogous to magnetic force in the gravitational case? And if it have its analogue, what is there to correspond with electric current? At first glance it might seem that the whole of the magnetic side of electromagnetism was absent in the gravitational analogy. But this is not true.

Thus, if u is the velocity of ρ , then ρu is the density of a current (or flux) of matter. It is analogous to a convective current of electrification. Also, when the matter ρ enters any region through its boundary, there is a simultaneous convergence of gravitational force into that region proportional to ρ . This is expressed by saying that if†

$$\mathbf{C} = \rho u - c\dot{e}, \quad \dots \quad (4)$$

then \mathbf{C} is a circuital flux. It is the analogue of Maxwell's true current; for although Maxwell did not include the convective term ρu , yet it would be against his principles to ignore it. Being a circuital flux, it is the curl of a vector, say

$$\text{curl } \mathbf{h} = \rho u - c\dot{e}. \quad \dots \quad (5)$$

This defines \mathbf{h} except as regards its divergence, which is arbitrary, and may be made zero. Then \mathbf{h} is the analogue of mag-

* "Thomson and Tait," Vol. I., Part II., § 549.

† Observe that it is $-e$ that is the analogue of electric force, and $-c\dot{e}$ of displacement current.

netic force, for it bears the same relation to flux of matter as magnetic force does to convective current. We have*

$$\begin{aligned} \mathbf{h} &= \text{curl pot } \mathbf{C}, \quad \dots \dots \dots (6) \\ &= \text{curl } \mathbf{A}, \end{aligned}$$

if $\mathbf{A} = \text{pot } \mathbf{C}$. But, since instantaneous action is here involved, we may equally well take

$$\mathbf{A} = \text{pot } \rho u, \quad \dots \dots \dots (7)$$

and its curl will be \mathbf{h} . Thus, whilst the ordinary potential P is the potential of the matter, the new potential \mathbf{A} is that of its flux.

Now if we multiply (5) by \mathbf{e} , we obtain

$$\mathbf{e} \text{ curl } \mathbf{h} = e\rho u - e\dot{c}\dot{e}, \quad \dots \dots \dots (8)$$

or, which is the same,†

$$\text{conv } \mathbf{Veh} = \mathbf{Fu} - \dot{\mathbf{U}}, \quad \dots \dots \dots (9)$$

if $U = \frac{1}{2}ce^2$. But $\dot{\mathbf{U}}$ represents the rate of exhaustion of potential energy, so $-\dot{\mathbf{U}}$ represents its rate of increase, whilst \mathbf{Fu} represents the activity of the force on ρ , increasing its kinetic energy. Consequently, the vector \mathbf{Veh} expresses the flux of gravitational energy. More strictly, any circuital flux whatever may be added. This \mathbf{Veh} is analogous to the electromagnetic \mathbf{VEH} found by Poynting and myself.‡ But there is a reversal of direction. Thus, comparing a single moving particle of matter with a similarly-moving electric charge, describe a sphere round each. Let the direction of motion be the axis, the positive pole being at the forward end. Then in the electrical case the magnetic force follows the lines of latitude with positive rotation about the axis, and the flux of energy coincides with the lines of longitude from the negative pole to the positive. But in the gravitational case, although \mathbf{h} still follows the lines of latitude positively, yet since the radial \mathbf{e} is directed to instead of from the centre, the flux of energy is along the lines of longitude from the positive pole to the

* See § 134, page 208.

† By the transformation (178), § 132.

‡ See § 70, equation (13). Take $\mathbf{e}_0 = 0 = \mathbf{h}_0$.

negative. This reversal arises from all matter being alike and attractive, whereas like electrifications repel one another.

The electromagnetic analogy may be pushed further. It is as incredible now as it was in Newton's time* that gravitative influence can be exerted without a medium; and, granting a medium, we may as well consider that it propagates in time, although immensely fast. Suppose, then, instead of instantaneous action, which involves†

$$\text{curl } \mathbf{e} = 0, \quad \dots \dots \dots (10)$$

we assert that the gravitational force \mathbf{e} in ether is propagated at a single finite speed v . This requires that

$$v^2 \nabla^2 \mathbf{e} = \ddot{\mathbf{e}}, \quad \dots \dots \dots (11)$$

for this‡ is the general characteristic of undissipated propagation at finite speed. Now,||

$$\nabla^2 = \nabla \text{ div} - \text{curl}^2,$$

so in space free from matter we have

$$-v^2 \text{curl}^2 \mathbf{e} = \ddot{\mathbf{e}}. \quad \dots \dots \dots (12)$$

But we also have, by (5),

$$-\text{curl } \mathbf{h} = \dot{c}\dot{\mathbf{e}}, \quad \dots \dots \dots (13)$$

away from matter. This gives a second value to $\ddot{\mathbf{e}}$, when we differentiate (13) to the time, say

$$\ddot{\mathbf{e}} = -\frac{1}{c} \text{curl } \dot{\mathbf{h}}. \quad \dots \dots \dots (14)$$

So, by (12) and (14), and remembering that we have already chosen \mathbf{h} circuital, we derive

$$cv^2 \text{curl } \mathbf{e} = \dot{\mathbf{h}}. \quad \dots \dots \dots (15)$$

* To Newton himself, as shown in his often-quoted letter to Bentley.

† This equation is equivalent to (2) above, and it implies that the gravitational force is *exactly* dependent on the configuration of the matter.

‡ As shown by Poisson, the value of a quantity at a given place and time, when controlled by this equation, depends solely on the state of things at distance vt and time t earlier. That is, disturbances travel at speed v . But it is much simpler to understand this property through plane waves.

See (193), § 132.

Or, if μ is a new constant, such that

$$\mu cv^2 = 1, \quad (16)$$

then (15) may be written in the form

$$\text{curl } \mathbf{e} = \mu \dot{\mathbf{h}}. \quad (17)$$

To sum up, the first circuital law (5), or

$$\text{curl } \mathbf{h} = \rho \mathbf{u} - c\dot{\mathbf{e}}, \quad (5) \text{ bis.}$$

leads to a second one, namely (17), if we introduce the hypothesis of propagation at finite speed.* This, of course, might be inferred from the electromagnetic case.

In order that the speed v should be not less than any value that may be settled upon as the least possible, we have merely to make μ be of the necessary smallness. The equation of activity becomes, instead of (9),†

$$\text{conv } \mathbf{Veh} = \mathbf{Fu} - \dot{\mathbf{U}} - \dot{\mathbf{T}}, \quad (18)$$

if $\mathbf{T} = \frac{1}{2}\mu\mathbf{h}^2$. The negative sign before the time-increase of this quantity points to exhaustion of energy, as before. If so, we should still represent the flux of energy by \mathbf{Veh} . But, of course, \mathbf{T} is an almost vanishing quantity when μ is small enough, or v big enough. Note that \mathbf{h} is not a negligible quantity, though the product $\mu\mathbf{h}$ is. Thus results will be sensibly as in the common theory of instantaneous action, although expressed in terms of wave-propagation. Results showing signs of wave-propagation would require an inordinately large velocity of matter through the ether. It may be worth while to point out that the lines of gravitational force connected with a particle of matter will no longer converge to it uniformly from all directions when the velocity v is finite, but will show a tendency to lateral concentration, though only to a sensible extent when the velocity of the matter is not an insensible fraction of v .

The gravitational-electromagnetic analogy may be further extended if we allow that the ether which supports and propagates the gravitational influence can have a translational‡

* These equations are analogous to (4) and (5), §35.

† This is analogous to (12), §70, with the impressed forces made zero.

‡ This does not exclude rotational motion, which is, in fact, a differential effect in a special kind of translational motion.

motion of its own, thus carrying about and distorting the lines of force. Making allowance for this convection of \mathbf{e} by the medium, with the concomitant convection of \mathbf{h} , requires us to turn the circuital laws (17), (18) to

$$\text{curl}(\mathbf{e} + \mu \mathbf{Vqh}) = \mu \dot{\mathbf{h}}, \quad (19)$$

$$\text{curl}(\mathbf{h} + c\mathbf{Veq}) = \rho \mathbf{u} - c\dot{\mathbf{e}}, \quad (20)$$

where \mathbf{q} is the velocity of the medium itself.*

It is needless to go into detail, because the matter may be regarded as a special and simplified case of my investigation of the forces in the electromagnetic field, with changed meanings of the symbols. It is sufficient to point out that the stress in the field now becomes prominent as a working agent. It is of two sorts, one depending upon \mathbf{e} and the other upon \mathbf{h} , analogous to the electric and magnetic stresses. The one depending upon \mathbf{h} is, of course, insignificant. The other consists of a pressure parallel to \mathbf{e} combined with a lateral tension all round it, both of magnitude $\frac{1}{2}c\mathbf{e}^2$. This was equivalently suggested by Maxwell. Thus two bodies which appear to attract are pushed together. The case of two large parallel material planes exhibits this in a marked manner, for \mathbf{e} is very small between them, and relatively large on their further sides.

But the above analogy, though interesting in its way, and serving to emphasise the non-necessity of the assumption of instantaneous or direct action of matter upon matter, does not enlighten us in the least about the ultimate nature of gravitational energy. It serves, in fact, to further illustrate the mystery. For it must be confessed that the exhaustion of potential energy from a universal medium is a very unintelligible and mysterious matter. When matter is infinitely widely separated, and the forces are least, the potential energy is at its greatest, and when the potential energy is most exhausted, the forces are most energetic!

Now there is a magnetic problem in which we have a kind of similarity of behaviour, viz., when currents in material circuits are allowed to attract one another. Let, for completeness, the initial state be one of infinitely wide separation of infinitely

* The additional terms are analogous to the motional electric and magnetic forces of §44 and §66, &c. See also (d), §132.

small filamentary currents in closed circuits. Then, on concentration to any other state, the work done by the attractive forces is represented by $\Sigma \frac{1}{2} \mu H^2$, where μ is the inductivity and H the magnetic force. This has its equivalent in the energy of motion of the circuits, or may be imagined to be so converted, or else wasted by friction, if we like. But, over and above this energy, the same amount, $\Sigma \frac{1}{2} \mu H^2$, represents the energy of the magnetic field, which can be got out of it in work. It was zero at the beginning. Now, as Lord Kelvin showed, this double work is accounted for by extra work in the batteries or other sources required to maintain the currents constant. (I have omitted reference to the waste of energy due to electrical resistance, to avoid complications.) In the gravitational case there is a partial analogy, but the matter is all along assumed to be incapable of variation, and not to require any supply of energy to keep it constant. If we asserted that $\frac{1}{2} ce^2$ was stored energy, then its double would be the work done per unit volume by letting bodies attract from infinity, without any apparent source. But it is merely the exhaustion of potential energy of unknown amount and distribution.*

Potential energy, when regarded merely as expressive of the work that can be done by forces depending upon configuration, does not admit of much argument. It is little more than a mathematical idea, for there is scarcely any physics in it. It explains nothing. But in the consideration of physics in general, it is scarcely possible to avoid the idea that potential energy should be capable of localisation equally as well as kinetic. That the potential energy may be itself ultimately kinetic is a separate question. Perhaps the best definition of the former is contained in these words:—Potential energy is energy that is not known to be kinetic. But, however this be, there is a practical distinction between them which it is found useful to carry out. Now, when energy can be distinctly localised, its flux can also be traced (subject to circuital indeterminateness, however). Also, this flux of energy forms a useful working idea when action at a distance is denied (even though the speed of transmission be infinitely great, or be

* It would appear that we must go to the ether to find the source of all energy.

assumed to be so). Any distinct and practical localisation of energy is therefore a useful step, wholly apart from the debatable question of the identity of energy advocated by Prof. Lodge.

From this point of view, then, we ought to localise gravitational energy as a preliminary to a better understanding of that mysterious agency. It cannot be said that the theory of the potential energy of gravitation exhausts the subject. The flux of gravitational energy in the form above given is, perhaps, somewhat more distinct, since it considers the flux only and the changes in the amount localised, without any statement of the gross amount. Perhaps the above analogy may be useful, and suggest something better.

Part II.

In the foregoing I partly assumed a knowledge on the part of the reader of my theory of convective currents of electrification* ("Electrical Papers," Vol. II., p. 495 and after), and only very briefly mentioned the modified law of the inverse squares which is involved, viz., with a lateral concentration of the lines of force. The remarks of the Editor† and of Prof. Lodge‡ on gravitational aberration, lead me to point out now some of the consequences of the modified law which arises when we assume that the ether is the working agent in gravitational effects, and that it propagates disturbances at speed v in the manner supposed in Part I. There is, so far as I can see at present, no aberrational effect, but only a slight alteration in the intensity of force in different directions round a moving body considered as an attractor.

Thus, take the case of a big Sun and small Earth, of masses S and E , at distance r apart. Let f be the unmodified force of S on E , thus

$$f = \frac{SE}{4\pi r^2 c}, \dots \dots \dots (1)$$

using rational units in order to harmonise with the electromagnetic laws when rationally expressed. Also, let F be the

* See also §§ 52 to 62, and §§ 163, 4.

† *The Electrician*, July 14, p. 277, and July 28, p. 340.

‡ *The Electrician*, July 28, p. 347.

modified force when the Sun is in motion at speed u through the ether. Then*

$$F = f \times \frac{1-s}{(1-s \sin^2 \theta)^{\frac{1}{2}}}, \quad \dots \quad (2)$$

where s is the small quantity u^2/v^2 , and θ is the angle between \mathbf{r} and the line of motion. ("Electrical Papers," Vol. II., pp. 495, 499).

Therefore, if the Sun is at rest, there is no disturbance of the Newtonian law, because its "field of force" is stationary. But if it has a motion through space, there is a slight weakening of the force in the line of motion, and a slight strengthening equatorially. The direction is still radial.

To show the size of the effect, let

$$\left. \begin{aligned} u &= 3 \times 10^7 \text{ centim. per sec.} \\ v &= 3 \times 10^{10} \text{ ,, ,, ,,} \end{aligned} \right\} \dots \quad (3).$$

This value of u is not very different from the speed attributed to fast stars, and the value of v is the speed of light itself.

So we have

$$s = \frac{u^2}{v^2} = \frac{1}{10^6}, \quad \dots \quad (4)$$

i.e., one-millionth. All perturbing forces of the first order are, therefore, of the order of magnitude of only one-millionth of the full force, even when the speed of propagation is as small as that of light.

The simplest case is when the common motion of the Sun and Earth is perpendicular to the plane of the orbit. Then $\theta = \frac{1}{2}\pi$ all round the orbit, and

$$F = f(1 + \frac{1}{2}s), \quad \dots \quad (5)$$

showing an increase in the force of attraction of S on E of one two-millionth part, without alteration of direction or variation in the orbit.†

* This is the case of steady motion. There is no simple formula when the motion is unsteady. Equation (2), above, is immediately derivable from equation (40), § 163.

† But Prof. Lodge tells me that our own particular Sun is considered to move only 10·9 miles per second. This is stupendously slow. The size of s is reduced to about 1/360 part of that in the text, and the same applies to the corrections depending upon it.

But when the common motion of the Sun and Earth is in their plane, θ varies from 0 to 2π in a revolution, so that the attraction on E, whilst towards the Sun's centre always, undergoes a periodic variation from

$$F = f(1 - s) \quad \dots \quad (6)$$

when $\theta = 0$, to

$$F = f(1 + \frac{1}{2}s), \quad \dots \quad (7)$$

when $\theta = \frac{1}{2}\pi$. The extreme variation is, therefore, $\frac{3}{2}sf$, according to the data used. The result is a slight change in the shape of the orbit.

But, to be consistent, having made v finite by certain suppositions, we should carry out the consequences more fully, and allow not merely for the change in the Newtonian law, as above, but for the force brought in by the finiteness of v which is analogous to the "electromagnetic force." This is very small truly, but so is the above change in the Newtonian law, and since they are of the same order of magnitude, we should also count the auxiliary force. Call it G . Then*

$$\mathbf{G} = \mathbf{F} \times \frac{xqu}{v^2} \times \mathbf{Vq}_1 \mathbf{Vr}_1 \mathbf{u}_1, \quad \dots \quad (8)$$

where F is as before, in (2) above, q is the actual speed of the Earth (not the same as u), and in the third vectorial factor, \mathbf{q}_1 , \mathbf{u}_1 , and \mathbf{r}_1 are unit vectors drawn parallel to the direction of the Earth's motion, of the Sun's motion, and from the Sun to the Earth. We see at once that the order of magnitude cannot be greater than that of the departure of F from f , before considered, because u and q will be of the same order, at least when u is big. As for x , it is simply a numerical factor, which cannot exceed 1, and is probably $\frac{2}{3}$.

The simplest case is when the motion of the Sun is perpendicular to the orbit of the Earth. Then

$$G = F \times xs \quad \dots \quad (9)$$

gives the tensor or size of the auxiliary force. It is radial, but outwards, so that the result is merely to reduce the size of

* See § 46, equation (9), and §§ 82 and 85. To apply to the present case, we have to note that G is the "electromagnetic force" of the Sun's "magnetic induction" on the Earth's "electric current," or rather, the analogue thereof; and also that the direction must be reversed, just as the ordinary "electric force" is taken reversed.

the previous correction, viz., the difference of F from f in the same motional circumstances.

But when the line of motion of the Sun is in the plane of the orbit, the case is much more complicated. The force G is neither constant (for the same distance) nor radial, except in four positions, viz., two in the line of motion of the Sun, when the auxiliary force vanishes, and two when $\theta = \pm \frac{1}{2}\pi$, when it is greatest. But this force is still in the plane of the orbit, which is an important thing, and is, moreover, periodic, so that the tangential component is as much one way as the other in a period.

All we need expect, then, so far as I can see from the above considerations, are small perturbations due to the variation of the force of gravity in different directions, and to the auxiliary force. Of course, there will be numerous minor perturbations.*

If variations of the force of the size considered above are too small to lead to observable perturbations of motion, then the striking conclusion is that the speed of gravity may even be the same as that of light. If they are observable, then, if existent, they should turn up, but if non-existent then the speed of gravity should be greater. Furthermore, it is to be observed that there may be other ways of expressing the propagation of gravity.

But I am mindful of the good old adage about the shoemaker and his last, and am, therefore, reluctant to make any more remarks about perturbations. The question of the ether in its gravitational aspect must be faced, however, and solved sooner or later, if it be possible. Perhaps, therefore, my suggestion may not be wholly useless.

* The solution for steady rectilinear motion has been employed. The justification thereof is the smallness of u/v and the large periodic time. If we allow for the small curvature of path of an attracting body, we shall introduce corrections of the second order of small quantities.